

## A simple eigenvalue sum-density relationship for atoms

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(Received May 30, 1984, revised August 21, 1984)

In recent years several attempts have been made [1-3] to obtain an energy-eigenvalue sum relationship for atoms and molecules. Such studies are useful in providing a sound theoretical basis for the successful Walsh and Woodward-Hoffmann rules which are based on the eigenvalue sum,  $E^s$ , in place of the total energy  $E$ . Unfortunately a general quantitative energy-eigenvalue sum relationship is not as yet available.

The purpose of this letter is to approach the problem from the one electron density,  $\rho(r)$ , point of view. We shall be concerned with a given isoelectronic series of atoms. We propose that within such a series

$$E^s \approx C_1 \langle \rho \rangle^{2/3} + C_2 \langle \rho \rangle^{1/3} \quad (1)$$

where  $E^s$  denotes the eigenvalue-sum  $\sum \nu_i \epsilon_i$  and the average density expectation value is given by  $\langle \rho \rangle = \int \rho^2 d\tau$ . Hyman et al. [4] have obtained a reasonably tight lower bound on  $E^{3/2}$  in terms of  $\langle \rho \rangle$  which is given by

$$\frac{\langle \rho \rangle}{E^{3/2}} \leq 0.205. \quad (2)$$

These authors noted that the numerical value of the ratio  $\langle \rho \rangle / E^{3/2}$  remains constant within an atomic  $N$ -electron series.

Based on the scaling arguments due to Szasz et al. [5] it has been shown [6] numerically that an excellent representation of the total electron-electron repulsion energy,  $V_{ee}$ , for atoms can be written as

$$V_{ee} = 0.3977 N^{4/3} \langle \rho \rangle^{1/3} \quad (3)$$

using Eqs. (2)–(3), respectively, in the form

$$E = C_1 \langle \rho \rangle^{2/3} \quad (4)$$

and

$$V_{ee} = C_2 \langle \rho \rangle^{1/3} \quad (5)$$

in the Hartree–Fock (HF) identity

$$E^s = E + V_{ee} \quad (6)$$

leads to Eq. (1).

In Table 1 we have listed the values of the fitting parameters  $C_1$  and  $C_2$  corresponding to the He, Li, and Be isoelectronic series along with the regression coefficient ‘ $r$ ’. All the fits are found to have excellent ‘ $r$ ’ values.

In Table 2 we have presented the values of  $\langle \rho \rangle^{1/3}$  and  $E^s$  based on the analytic Hartree–Fock wave functions [7] and the values of  $E^s$  obtained from Eq. (1). These results show that Eq. (1) provides a good quantitative relationship for  $E^s$  in terms of  $\langle \rho \rangle$ . We also note that the quantity  $2\langle \rho \rangle^{1/3}$  closely approximates the value of  $Z$ , the atomic number, and it is therefore tempting [8] to represent  $E^s$  and the total HF energy  $E$  in terms of a polynomial in  $\langle \rho \rangle^{1/3}$ . A polynomial fit

**Table 1.** Fitting parameters  $C_1$  and  $C_2$  in Eq. (1) for the He, Li, and Be isoelectronic series respectively

Series	Fitting parameters		Regression coefficient
	$-C_1$	$C_2$	
He	3.3975	1.1063	0.9999
Li	3.5160	1.6712	0.9999
Be	3.5972	2.3937	0.9999

**Table 2.** The values of  $\langle \rho \rangle^{1/3}$  and  $E^s(H.F.)$  based on the Hartree–Fock wave functions (Ref. 4 and Ref. 7) for some members of the He, Li, and Be series.  $E^s$  values obtained from Eq. (1) are given in the last column

Series	$\langle \rho \rangle^{1/3}$	$E^s(H.F.)$	$E^s$ (Fitted)
He	0.9153	1.836	1.831
Li <sup>+</sup>	1.4551	5.585	5.582
N <sup>5+</sup>	3.6225	40.584	40.579
Li	1.4642	5.152	5.091
Be <sup>+</sup>	2.0210	10.943	10.984
O <sup>5+</sup>	4.2572	56.723	56.609
Be	2.0325	10.076	9.994
B <sup>+</sup>	2.6060	18.119	18.192
O <sup>4+</sup>	4.3353	57.366	57.232

Table 3. The fitting parameters given in Eq. (7) for He, Li, and Be-series corresponding to the total Hartree-Fock energy,  $E^{HF}$ , and the Hartree-Fock eigenvalue sum,  $E^S$ , respectively. The number( $n$ ) denotes  $\times 10^n$ . All values are in a.u.

Series	Fitting parameters					Average % error
	A	B	C	D	E	
He $E^{HF}$	-3.406	-0.940 (-2)	-0.101 (-1)	-0.192 (-1)	0.261 (-1)	0.3 (-3)
He $E^S$	-3.406	1.152	-0.583 (-1)	0.999 (-2)	0.937 (-2)	0.6 (-3)
Li $E^{HF}$	-3.590	0.243	-0.189	0.253	-0.161	0.8 (-3)
Li $E^S$	-3.590	0.243	-0.189	0.253	-0.161	0.8 (-3)
Be $E^{HF}$	-3.738	0.519	-0.381	0.696	-0.613	0.1 (-3)
Be $E^S$	-3.729	3.111	-0.560	-0.921	0.705 (-1)	0.1 (-3)

of  $E^s$  (and  $E^{HF}$ ) according to

$$E^s = A\langle\rho\rangle^{2/3} + B\langle\rho\rangle^{1/3} + C + D\langle\rho\rangle^{-1/3} + E\langle\rho\rangle^{-2/3} \quad (7)$$

is given in Table 3 for the He, Li and Be series respectively. Such polynomials can be used to obtain the estimates of  $E^s$  from the experimentally determined  $\langle\rho\rangle$  values.

We wish to point out that Eq. (1) is only a semiempirical result in the sense that although its form can be arrived at in a purely theoretical manner, the actual constants  $C_1$  and  $C_2$  have been obtained using the available  $HF$  data. In the light of the existence of a unique energy density functional for the ground state of an  $N$ -electronic system [9] and Eq. (1), we conclude that the energy-eigenvalue sum relationship can be understood in terms of the single particle density.

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